

# Random Convex Hulls: A Variance Revisited

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## Abstract

An exact expression is determined for the asymptotic constant  $c_2$  in the limit theorem of GROENEBOOM [4] (1988) which states that the number of vertices of the convex hull of a uniform sample of  $n$  random points from a circular disk satisfies a central limit theorem as  $n \rightarrow \infty$  with asymptotic variance  $2\pi c_2 n^{1/3}$ .

## 1 Introduction

The convex hull of random point sets that arises in the statistical analysis of multivariate data, as explained elsewhere a great number of times, forms a vital and widely studied subject in stochastic geometry. An alternative point of view looks at the convex hull of a random point set as a graph on random points in which the vertices of the convex hull are connected. A multitude of questions, especially in three and higher dimensions and in the non-uniform case, remain open.

Consider the convex hull  $K_n$  of  $n$  independent and uniformly distributed points, drawn from a fixed circular disk  $K$  in  $\mathbf{R}^2$ . Let  $N_n$  and  $D_n$  denote the number of vertices of the random polygon  $K_n$  and the area of  $K \setminus K_n$ , the area in  $K$  outside the convex hull  $K_n$ , respectively. GROENEBOOM [4] in 1988 made substantial progress in our understanding of  $N_n$ . He introduced an elegant method to derive its asymptotic  $k$ -th moments for  $k \geq 1$  and proved the central limit theorem that as  $n \rightarrow \infty$ ,

$$\frac{N_n - 2\pi c_1 n^{1/3}}{\{2\pi c_2 n^{1/3}\}^{1/2}} \xrightarrow{d} Z \sim \mathcal{N}(0, 1), \quad (1.1)$$

where “ $\xrightarrow{d}$ ” denotes convergence in distribution,  $\mathcal{N}(0, 1)$  stands for the standard normal distribution,  $c_1 = (3\pi/2)^{-1/3} \Gamma(\frac{5}{3}) \approx 0.5384576135\dots$  and  $c_2$  is expressed in terms of two rather complicated double integrals whose approximate values are obtained numerically in his paper. (We will see shortly that his stated numerical value 0.33503... for  $c_2$  is incorrect.)

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Another result in [4] says that, when  $K$  is replaced by a fixed convex polygon with  $r \geq 3$  vertices in the plane, then  $N_n$  as  $n \rightarrow \infty$  satisfies a central limit theorem with asymptotic mean  $(2/3)r \log n$  and variance  $(10/27)r \log n$  (HUETER [6] proves a central limit theorem for  $N_n$  when the sample is normal rather than uniform). While the exact value of  $c_1$  (in the case of a circular disk) goes back to the work of RÉNYI AND SULANKE [7] in 1963, to the best of the authors' knowledge, it is an open problem to determine the exact value for  $c_2$ . The purpose of this note is to do so. BUCHTA [2] points out at the conclusion of Section 5.2 that finding an exact expression for  $c_2$  may also identify the variance in a related question: HSING [5] in 1994 demonstrated that

$$\lim_{n \rightarrow \infty} n^{5/3} \text{Var}(D_n) = \sigma^2,$$

and as  $n \rightarrow \infty$ ,

$$n^{5/6} (D_n - E(D_n)) \xrightarrow{d} \tilde{Z} \sim \mathcal{N}(0, \sigma^2) \quad (1.2)$$

for a finite nonnegative number  $\sigma^2$  (BUCHTA refers to  $\sigma^2$  as Hsing's constant). The parameter  $\sigma^2$  is in terms of an integral. How to evaluate this integral is not obvious nor is the strict inequality  $\sigma^2 > 0$  apparent. Let  $\text{area}(K)$  denote the area of the disk  $K$ . BUCHTA [2] explains that, if  $\text{Var}(N_n) \sim 2\pi c_2 n^{1/3}$  as  $n \rightarrow \infty$ , then

$$\sigma^2 = 2\pi \left(\frac{1}{3} c_1 + c_2\right) (\text{area}(K))^2 = \theta (\text{area}(K))^2. \quad (1.3)$$

We remark that it is stated (without proof) in GROENEBOOM [4], page 328, that  $\text{Var}(N_n) \sim 2\pi c_2 n^{1/3}$  as  $n \rightarrow \infty$ , in other words, the variance of the asymptotic normal distribution of  $N_n$  equals the asymptotic variance of  $N_n$ . Under the assumption that the relationship in (1.3) between  $\sigma^2$  and  $c_2$  holds, applying the Efron identity (EFRON [3], 1965)

$$\frac{ED_n}{\text{area}(K)} = \frac{EN_{n+1}}{n+1}$$

in combination with the result of RÉNYI AND SULANKE on  $EN_n$ , we may restate the limit theorem in (1.2) as

$$\frac{(\text{area}(K))^{-1} D_n - 2\pi c_1 n^{-2/3}}{\{2\pi(\frac{1}{3} c_1 + c_2) n^{-5/3}\}^{1/2}} \xrightarrow{d} Z \sim \mathcal{N}(0, 1) \quad (1.4)$$

as  $n \rightarrow \infty$ . Our contribution is the following

**THEOREM 1** *The asymptotic constant  $c_2$  is given by*

$$c_2 = \frac{16\pi^2 \Gamma(\frac{2}{3})^{-3} - 57}{27} c_1 \approx 0.1316029298\dots, \quad (1.5)$$

where  $c_1 = (3\pi/2)^{-1/3} \Gamma(\frac{5}{3}) \approx 0.5384576135\dots$ , and the constant  $\theta$  is

$$\theta = 2\pi \left(\frac{1}{3} c_1 + c_2\right) = \frac{32\pi(\pi^2 \Gamma(\frac{2}{3})^{-3} - 3)}{27} c_1 \approx 1.9546285835\dots \quad (1.6)$$

**Remarks.** (1) We see from (1.5) that the limiting ratio  $c_2/c_1$  of the variance and the expectation of the asymptotic normal distribution of  $N_n$  as  $n \rightarrow \infty$  is 0.2444072226..., which is smaller than  $5/9$ , the analogous ratio in the case when a convex polygon with  $r \geq 3$  vertices in the plane replaces the disk  $K$ .

(2) The discrepancy in the numerical value for  $c_2$  in [4] might be explainable by the factor 2 in front of the double integral in the formula after display (3.35) not having been included in the numerical evaluation.

## 2 Proof of Theorem 1

In order to determine an exact expression for  $c_2$  in Theorem 3.4 in [4], it suffices to do so for  $c_2$  in Theorem 3.3(i) in [4] by deriving an asymptotic expression for  $\text{Var}(N(a))$  as  $a \rightarrow \infty$ . For convenience, we recall a couple of definitions from [4] that lead up to the one of  $N(a)$ .

Let  $\{T(a) = (R(a), S(a)) : a \in \mathbf{R}\}$  denote the stationary jump process that arises from the jump process  $\{Z(a) = (X(a), Y(a)) : a \in \mathbf{R}\}$  (see Definition 3.3) upon the transformations in (3.18) and (3.19) in [4], p. 356. The process  $\{Z(a) : a \in \mathbf{R}\}$  relies on the approximation of the sample point process of size  $n$  from the unit disk by a Poisson point process on  $\mathbf{R}^2$  with intensity  $(n/\pi) \times$  Lebesgue measure (see Lemma 3.2). With these defined,  $N(a)$  denotes the number of jumps of the process  $\{T(c) : c \geq 0\}$  on the time interval  $[0, a]$  (p. 357, [4]). We refer the interested reader to [4] for the full development that concludes with the limit law for  $N_n$ .

Lemma 3.7 lies at the core of the main results in [4]. Since the former result has a few misprints at crucial places in such a fashion that it is rather challenging to follow the exposition, starting on page 360 and running into page 361 of [4], we give a corrected version of that lemma: For  $u, v > 0$ , define

$$\begin{aligned} A_1(u, v) &= \sqrt{u} + \frac{1}{2}(u - v + 1) \\ A_2(u, v) &= \sqrt{v} + \frac{1}{2}(v - u + 1) \\ A_3(u, v) &= \frac{2}{3}(u^{3/2} + v^{3/2}) + \frac{1}{4}(u - v)^2 + \frac{1}{2}(u + v) - \frac{1}{12} \end{aligned} \quad (2.7)$$

and let  $\mathcal{K}$  be the region  $\mathcal{K} = \{(u, v) \in \mathbf{R}_+^2 : |1 - v^{1/2}| < u^{1/2} < 1 + v^{1/2}\}$ . Then for  $c > 0$ ,

$$\begin{aligned} E[(S(0)^{1/2} + R(0))^2 \cdot ((S(c)^{1/2} - R(c))^2) &= \\ &= \frac{c^{10}}{36\pi^2} \iint_{\mathcal{K}} A_1(u, v)^3 A_2(u, v)^3 e^{-A_3(u, v)c^3/(2\pi)} du dv \\ &+ \frac{c^7}{4\pi} \iint_{\mathcal{K}} A_1(u, v)^2 A_2(u, v)^2 e^{-A_3(u, v)c^3/(2\pi)} du dv \\ &+ \frac{16}{9\pi^2} \iint_{c - \sqrt{v} > \sqrt{u}} u^{3/2} v^{3/2} e^{-2(u^{3/2} + v^{3/2})/(3\pi)} du dv. \end{aligned} \quad (2.8)$$

(The factor  $1/4$  in  $A_3(u, v)$  in (2.7) was  $1/2$  and the ratio  $16/(9\pi^2)$  in (2.8) was  $4/(9\pi^2)$  in [4].) Proceeding along the arguments relied on after display (3.35) in [4], as  $a \rightarrow \infty$ ,

$$\begin{aligned}
\text{Var}(N(a)) &\sim c_1 a + 2 \int_0^a db \int_0^b \frac{1}{(2\pi)^2} E[(S(0)^{1/2} + R(0))^2 (S(c)^{1/2} - R(c))^2] dc - (c_1 a)^2 \\
&= c_1 a + \frac{1}{2\pi^2} \int_0^a db \int_0^b \left[ \frac{c^{10}}{36\pi^2} \iint_{\mathcal{K}} A_1(u, v)^3 A_2(u, v)^3 e^{-A_3(u, v)c^3/(2\pi)} du dv \right. \\
&\quad + \frac{c^7}{4\pi} \iint_{\mathcal{K}} A_1(u, v)^2 A_2(u, v)^2 e^{-A_3(u, v)c^3/(2\pi)} du dv \\
&\quad + \frac{16}{9\pi^2} \left\{ \iint_{c-\sqrt{v} > \sqrt{u}} u^{3/2} v^{3/2} e^{-2(u^{3/2}+v^{3/2})/(3\pi)} du dv \right. \\
&\quad \left. - \int_0^\infty u^{3/2} e^{-2u^{3/2}/(3\pi)} du \int_0^\infty v^{3/2} e^{-2v^{3/2}/(3\pi)} dv \right\} \Big] dc \\
&= c_1 a + \frac{1}{2\pi^2} \int_0^a db \left[ \frac{1}{36\pi^2} \iint_{\mathcal{K}} A_1(u, v)^3 A_2(u, v)^3 B_1(u, v) du dv \right. \\
&\quad + \frac{1}{4\pi} \iint_{\mathcal{K}} A_1(u, v)^2 A_2(u, v)^2 B_2(u, v) du dv \\
&\quad \left. - \frac{16}{9\pi^2} \iint_{c-\sqrt{v} < \sqrt{u}} u^{3/2} v^{3/2} e^{-2(u^{3/2}+v^{3/2})/(3\pi)} du dv \right] \\
&\sim a \left[ c_1 + \frac{1}{2\pi^2} \left\{ \frac{1}{36\pi^2} \frac{1}{3} (2\pi)^{11/3} \Gamma(11/3) I_1 + \frac{1}{4\pi} \frac{1}{3} (2\pi)^{8/3} \Gamma(8/3) I_2 - \frac{16}{9\pi^2} I_3 \right\} \right], \tag{2.9}
\end{aligned}$$

where  $B_1(u, v)$  and  $B_2(u, v)$  are given at the bottom of page 361 in [4],  $I_1$  and  $I_2$  are the two double integrals

$$\begin{aligned}
I_1 &= \iint_{\mathcal{K}} A_1(u, v)^3 A_2(u, v)^3 A_3(u, v)^{-11/3} du dv, \\
I_2 &= \iint_{\mathcal{K}} A_1(u, v)^2 A_2(u, v)^2 A_3(u, v)^{-8/3} du dv,
\end{aligned}$$

and the integral  $I_3$  is evaluated in [4], p. 362, and comes out as  $I_3 = \frac{8}{9}(3\pi/2)^{11/3} \Gamma(5/3)$ .

Let us now turn to the evaluation of  $I_1$  and  $I_2$ . For this purpose, we make the substitution  $u = (x+y)^2$  and  $v = (x-y)^2$ , equivalently,  $x = (u^{1/2} + v^{1/2})/2$  and  $y = (u^{1/2} - v^{1/2})/2$ . This transforms the region  $\mathcal{K}$  into the semi-infinite strip  $\tilde{\mathcal{K}} = \{(x, y) : 1/2 \leq x < \infty, -1/2 \leq y \leq 1/2\}$ . Also, we obtain  $du dv = 8(x^2 - y^2) dx dy$ , while  $A_1(u, v)A_2(u, v)$  and  $A_3(u, v)$  become  $A_1(x, y)A_2(x, y) = (2x+1)^2(1-4y^2)/4$  and  $A_3(x, y) = (2x+1)^2(\frac{x}{3} + y^2 - \frac{1}{12})$ . Consequently,

$$\begin{aligned}
I_1 &= \frac{1}{8} \iint_{\tilde{\mathcal{K}}} (x^2 - y^2) (2x+1)^{-4/3} (1-4y^2)^3 \left(\frac{x}{3} + y^2 - \frac{1}{12}\right)^{-11/3} dx dy \\
I_2 &= \frac{1}{2} \iint_{\tilde{\mathcal{K}}} (x^2 - y^2) (2x+1)^{-4/3} (1-4y^2)^2 \left(\frac{x}{3} + y^2 - \frac{1}{12}\right)^{-8/3} dx dy. \tag{2.10}
\end{aligned}$$

Some computations that we report in the Appendix provide

$$I_1 = -(9/8)6^{2/3} + (27/5)2^{-1/3} 3^{7/6} I_4 = 4.1603038934\dots, \tag{2.11}$$

where  $I_4 = \int_{4^{-1/3}}^1 (4t^3 - 1)^{-1/2} dt$ . A glimpse at formulas (18.13.17) and (18.13.24) of ABRAMOWITZ AND STEGUN [1] reveals that, alternatively,  $I_4 = \Gamma(1/3)^3/(12\pi)$ . Similarly, in a couple of steps, one derives an expression for  $I_2$ , which simplifies to

$$I_2 = 6^{5/3}/5 = 3.9623126986\dots \quad (2.12)$$

Furthermore, one shows that  $I_1$  and  $I_2$  are related by

$$I_1 = \frac{8\pi^2 \Gamma(2/3)^{-3} - 15}{16} I_2. \quad (2.13)$$

Finally, plugging  $I_1$ ,  $I_2$ , and  $I_3$  into the expression for  $\text{Var}(N(a))$  in (2.9), appealing to the identity  $\Gamma(w+1) = w\Gamma(w)$  repeatedly as necessary, and recalling the formula for  $c_1$ , after a sequence of elementary steps, we arrive at, as  $a \rightarrow \infty$ ,

$$\text{Var}(N(a)) \sim \frac{a c_1}{27} [16\pi^2 \Gamma(2/3)^{-3} - 57].$$

This completes the proof of the statement in (1.5). Since the assertion in (1.6) is readily verified, this finishes our proof of Theorem 1.

## A Appendix: Identity in (2.11)

We use the shorthand notation  $\alpha = y^2 - 1/12$ ,  $\beta = 1 - 6\alpha$  and  $\gamma = x + 3\alpha$  and write

$$\begin{aligned} g_1(x, y) &= (\gamma/3)^{1/3} (2x+1)^{2/3} \{(-243/8)\beta^{-2}\gamma^{-3} + (1701/10)\beta^{-3}\gamma^{-2} \\ &\quad - (9963/10)\beta^{-4}\gamma^{-1} - 1944\beta^{-4}(2x+1)^{-1}\} \\ g_2(x, y) &= (\gamma/3)^{1/3} (2x+1)^{2/3} \{(-2187/8)\alpha^2\beta^{-2}\gamma^{-3} + (729/10)\alpha(4-3\alpha)\beta^{-3}\gamma^{-2} \\ &\quad + (243/10)(-5-36\alpha+27\alpha^2)\beta^{-4}\gamma^{-1} - 486\beta^{-4}(2x+1)^{-1}\} \\ g_3(y) &= -(27/20)6^{2/3}(7/4+y^2) \\ g_4(y) &= -54 \cdot 3^{2/3}(1+12y^2)^{-8/3}(1/5+4y^2/5+144y^4/5+64y^6) \\ G_3(y) &= -(9/20)6^{2/3}y(21/4+y^2) \\ G_4(y) &= -(27/5)3^{2/3}y(1+12y^2)^{-5/3}(1+4y^2+32y^4). \end{aligned}$$

Integrating with respect to  $x$ , we obtain

$$\begin{aligned} I_1 &= \frac{1}{8} \int_{-1/2}^{1/2} [g_2(x, y) - y^2 g_1(x, y) \Big|_{1/2}^{\infty}] \cdot (1-4y^2)^3 dy \\ &= \int_{-1/2}^{1/2} (g_3(y) - g_4(y)) dy \\ &= (G_3(y) - G_4(y)) \Big|_{-1/2}^{1/2} + (27/5)3^{2/3} \int_{-1/2}^{1/2} (1+12y^2)^{-2/3} dy \\ &= -(9/8)6^{2/3} + (27/5)3^{2/3}2^{-1/3}3^{1/2} \int_{4^{-1/3}}^1 (4t^3-1)^{-1/2} dt \quad (3.1) \\ &= -(9/8)6^{2/3} + (27/5)2^{-1/3}3^{7/6} I_4, \end{aligned}$$

where the substitution  $t = 4^{-1/3} (1 + 12y^2)^{1/3}$  provides the Weierstrass elliptic integral in (3.1).

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